

# Kac's Program for the Landau Equation

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# Description of the Physical World

**Basic Question:** How can we describe a complicated system of a vast number of identical interacting particles, such as a rarefied gas or a plasma?

**Microscopic scale:** state vector  $(x_1, v_1, \dots, x_N, v_N) \in (\mathbb{R}^3 \times \mathbb{R}^3)^N$  governed by Newton's second law of motion. Complete and detailed, but impossible to track and analyze mathematically.

**Macroscopic scale:** spatial observables such as the mass  $\rho(t, x)$ , the velocity  $u(t, x)$  and the temperature  $T(t, x)$  governed by fluid equations. Losing information about the velocity distribution.

**Mesoscopic scale:** one-particle density function  $f(t, x, v)$  governed by kinetic equations. Intermediate between the above two scales.

# Foundations of Kinetic Theory

**Maxwell** (1867): started the study of dynamics of gases; wrote down the first version of the Boltzmann equation and discovered the Maxwellian equilibrium.

**Boltzmann** (1872): wrote down the explicit formula of the Boltzmann equation; established the famous  $H$ -theorem and explained the convergence to the Maxwellian equilibrium.

**Landau** (1936): wrote down the Landau equation as an alternative to the Boltzmann equation when the charged particles perform Coulomb collisions in a plasma, by assuming that the grazing collisions prevail.

**Hilbert** (1900): asked for a rigorous derivation from the microscopic Newton's law to the mesoscopic kinetic equations and then to the macroscopic fluid equations in his sixth problem at the ICM.

The essentially difficult first part remains open for the Landau equation.

Bobylev-Pulvirenti-Saffirio 2013 CMP: a consistency result.

Lanford 1975, Deng-Hani-Ma 2024, 2025: from hard-sphere to Boltzmann.

# Kac's Program

**Kac** (1956): introduced his seminal program to derive *spatially homogeneous* kinetic equations from many-particle systems.

**Kac's model**: starting from a conservative stochastic model, which simplifies the dynamics to a random Markovian process but also preserves some key features.

**Boltzmann collisional operator**:  $B(|z|, \cos \theta) = \Gamma(|z|) \cdot b(\cos \theta)$ .

**Initial state space**:  $N$  i.i.d. velocities  $V_1, \dots, V_N$  with initial law  $f_0$  in  $\mathbb{R}^3$ .

**Collision rate**: for each pair of indices  $(i, j)$ , choose a random collision time  $T_{ij}$  according to the exponential distribution with parameter  $\Gamma(|V_i - V_j|)$ . Let  $T = \min T_{ij}$  be the collision time and change the corresponding pair of velocities.

**Post-collisional state space**:

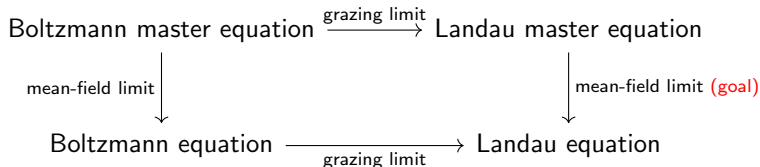
$$V'_i = \frac{V_i + V_j}{2} + \frac{|V_i - V_j|}{2} \sigma, \quad V'_j = \frac{V_i + V_j}{2} - \frac{|V_i - V_j|}{2} \sigma,$$

where  $\sigma \in \mathbb{S}^2$  is drawn according to the law  $b(\cos \theta_{ij})$ .  $\theta_{ij}$  is the deviation angle between the pre-collisional and post-collisional velocities.

# Kac's Program for the Landau Equation

The Landau equation is formally obtained by passing the Boltzmann equation to the *grazing limit*.

We can also pass to the *grazing limit* for the Boltzmann master equation of Kac's stochastic model to obtain the Landau master equation.



# Landau Equation and Master Equation

**Landau Equation:** integral-differential form

$$\partial_t f = \nabla_v \cdot \int_{\mathbb{R}^3} a(v-w) \cdot (\nabla_v - \nabla_w) f(v) f(w) dw,$$

or more commonly used divergence form

$$\partial_t f = \nabla \cdot [(a * f) \nabla f - (b * f) f].$$

Here  $a(z) = |z|^\gamma (|z|^2 \text{Id} - z \otimes z)$  and  $b(z) = \nabla \cdot a(z) = -2|z|^\gamma z$ .

The parameter  $\gamma \in [-3, 1]$  represents the regularity of the interaction forces.  
 $\gamma = -3$  corresponds to the most important Coulomb interaction case.

**Landau Master Equation:**

$$\partial_t F_N = \frac{1}{N} \sum_{i < j}^N (\nabla_{v_i} - \nabla_{v_j}) \cdot \left( a(v_i - v_j) \cdot (\nabla_{v_i} - \nabla_{v_j}) F_N \right).$$

We take  $\gamma = -3$  in this talk.

# Kac's Program for the Landau Equation

## Theorem (F.-Z.Wang (2025))

Let  $f_0 \in L^1 \cap L^\infty(\mathbb{R}^3)$  be a probability density. Let  $F_N \in L^\infty([0, T]; L^1 \cap L^\infty(\mathbb{R}^{3N}))$  be the unique bounded weak solution to the Landau master equation with initial data  $f_0^{\otimes N}$ . Let  $f \in C^1((0, \infty); \mathcal{S}(\mathbb{R}^3)) \cap L^\infty([0, \infty); L^1 \cap L^\infty(\mathbb{R}^3))$  be the unique bounded smooth solution to the Landau equation with initial data  $f_0$ . Assume that  $f_0$  has finite weighted Fisher information and high-order  $L^1$  moment:

$$\int_{\mathbb{R}^3} \langle v \rangle^3 |\nabla \log f_0|^2 f_0 \, dv < \infty, \quad \int |v|^m f_0 \, dv < \infty \text{ for some } m > 6.$$

Then for any  $T > 0$ , we have propagation of chaos: for any  $k \geq 1$ , the  $k$ -marginal  $F_{N,k}$  converges weakly to  $f^{\otimes k}$  as  $N \rightarrow \infty$  on  $[0, T] \times \mathbb{R}^{3k}$ :

$$\int_{\mathbb{R}^{3k}} (F_{N,k} - f^{\otimes k}) \varphi^{\otimes k} \, dv_1 \cdots dv_k \rightarrow 0$$

as  $N \rightarrow \infty$ , for any  $\varphi \in C_c^\infty(\mathbb{R}^3)$  and any  $t \in [0, T]$ . Here  $\langle v \rangle = \sqrt{1 + |v|^2}$ .

# Further Discussion

- 1 Quantitative propagation of chaos for the Landau master equation;  
Convergence in better sense.
- 2 Similar results for the Boltzmann equation or the Vlasov-Poisson equation?
- 3 Central limit theorem/ Gaussian fluctuation/ Large deviation principles for kinetic equations.
- 4 Derivation of spatially inhomogeneous kinetic equations/ from Newton dynamics under low-density limit or weak-coupling limit.



# References



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# Thank you!